

# Physics 2D, Winter 2005

## Week 4 Exercise Solutions

$$4] a] e_{\text{TOT}} = \sigma T^4 = 5.67 \times 10^{-8} \frac{\text{W}}{\text{m}^2 \text{K}^4} (3000 \text{K})^4$$
$$= \boxed{4.6 \times 10^6 \text{ W/m}^2}$$

$$b] 75 \text{ W} = e_{\text{TOT}} \cdot \text{Area}$$
$$\Rightarrow \text{Area} = 1.63 \times 10^{-5} \text{ m}^2 = \boxed{16.3 \text{ mm}^2}$$

$$10] 1 \text{ photon has } E = hf = (6.626 \times 10^{-34} \text{ J}\cdot\text{s})(94 \text{ MHz})$$
$$= 6.23 \times 10^{-26} \text{ J}$$

$$100 \text{ kW} = \frac{10^5 \text{ J}}{\text{s}}, \text{ so this means } \frac{10^5}{6.23 \times 10^{-26}} = \boxed{1.6 \times 10^{30} \text{ photons/sec}}$$

$$13] \text{ So } K_{\text{max}} = 2.92 \text{ eV.}$$

$$\text{Thus } 2.92 \text{ eV} = hf - \phi$$

$$\text{But } hf = \frac{hc}{\lambda} = \frac{1240 \text{ eV}\cdot\text{nm}}{250 \text{ nm}} = 4.96 \text{ eV}$$

$$\text{So } \phi = (4.96 - 2.92 \text{ eV}) = \boxed{2.04 \text{ eV}}$$

$$17 \text{ a) } E_{\text{light}} = \frac{hc}{\lambda} = 4.13 \text{ eV}$$

So Lithium & beryllium will eject electrons, since  
 $\phi < E_{\text{light}}$

$$b) K_{\text{max}} = hf - \phi = \begin{cases} 1.83 \text{ eV} & \text{for Lithium} \\ 0.23 \text{ eV} & \text{for beryllium.} \end{cases}$$

$$20) K_{\text{max}} = \frac{hc}{\lambda} - \phi, \text{ so for the first wavelength,}$$

$$1 \text{ eV} = \frac{hc}{\lambda} - \phi. \text{ The second wavelength is } \frac{\lambda}{2}, \text{ and so}$$

$$4 \text{ eV} = \frac{2hc}{\lambda} - \phi. \text{ We want to solve for } \phi, \text{ so just}$$

divide the 2<sup>nd</sup> eqn by 2 and subtract:

$$1 \text{ eV} = \frac{hc}{\lambda} - \phi$$

$$- (2 \text{ eV} = \frac{hc}{\lambda} - \frac{\phi}{2})$$

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$$-1 \text{ eV} = -\frac{\phi}{2} \Rightarrow$$

$$\boxed{\phi = 2 \text{ eV}}$$

22 | Need to find  $K_{\max}$ , so need  $v$  for the electrons.

For  $e^-$ 's in a magnetic field, we know  $F = \boxed{qvB = \frac{mv^2}{r}}$

(it's safe to use non-relativistic mechanics here, since these are not very energetic electrons).

$$\text{So } v = \frac{qrB}{m}, \text{ and } K = \frac{1}{2}mv^2 = \frac{1}{2} \frac{q^2 r^2 B^2}{m}$$

↑  
again, nonrelativistic is OK.

25 | a)  $\lambda' - \lambda_0 = \frac{h}{mec} (1 - \cos \theta)$

$$= \frac{h}{mec} (1 - \cos 30^\circ)$$

$$= 3.25 \times 10^{-13} \text{ m} = \boxed{3.25 \times 10^{-4} \text{ nm}}$$

b)  $\lambda' = \lambda_0 + \frac{h}{mec} (1 - \cos \theta)$ .

And  $\lambda_0 = \frac{hc}{E} = \frac{1240 \text{ eV nm}}{300 \text{ keV}} = 4.13 \times 10^{-3} \text{ nm}$ .

So  $\lambda' = 4.46 \times 10^{-3} \text{ nm} \Rightarrow E' = \frac{hc}{\lambda} = \frac{1240 \text{ eV nm}}{4.46 \times 10^{-3} \text{ nm}} = \boxed{278 \text{ keV}}$

c)  $KE_{e^-} = E' - E = \boxed{22 \text{ keV}}$

(That's just the kinetic energy - you could include rest energy too if you wanted).

27 | The setup:

Before  $\rightarrow \cdot e^-$

After  $e^- \rightarrow$

Energy |  $E_\gamma + mc^2 = \cancel{E_\gamma} E_e$

Momentum |  $p_\gamma = p_e$

So since  $E_\gamma = p_\gamma c$  and  $E_e = \sqrt{p_e^2 c^2 + m^2 c^4}$   
we can combine these to say

$$\boxed{p c + m c^2 = \sqrt{p^2 c^2 + m^2 c^4}} \quad (p = p_e = p_\gamma)$$

Square both sides:  $p^2 c^2 + 2 p m c^3 + m^2 c^4 = p^2 c^2 + m^2 c^4$

$$\Rightarrow \boxed{2 p m c^3 = 0}$$

This means  $p = 0$ , which says there's no photon to begin with. So the only solution is to have a free electron just hanging out.

30] We know  $\lambda' = \lambda_0 + \frac{h}{m_e c} (1 - \cos \theta)$ .

For max energy transfer,  $\lambda$  should increase (i.e. the energy of the photon should decrease) as much as possible

So  $\theta = 180^\circ \Rightarrow \boxed{\lambda' = \lambda_0 + \frac{2h}{m_e c}}$

Now,  $E = \frac{hc}{\lambda}$  for a photon.

Let's call  $E' = \frac{hc}{\lambda'}$ ,  $E_0 = \frac{hc}{\lambda_0}$ . So then

$$E' = \frac{hc}{\lambda'} = \frac{hc}{\left(\lambda_0 + \frac{2h}{m_e c}\right)}$$

Let's assume  $\frac{2h}{m_e c} \ll \lambda_0$  and check this later.

Then  $\frac{1}{\lambda_0 + \frac{2h}{m_e c}} = \frac{1}{\lambda_0 \left(1 + \frac{2h}{m_e c \lambda_0}\right)} \approx \frac{1}{\lambda_0} \left(1 - \frac{2h}{m_e c \lambda_0}\right)$

So  $\boxed{E' = E_0 - \frac{2h^2 c^2}{m_e c^2 \lambda_0^2}}$

(I multiplied top & bottom of the fraction by  $c$ )

This makes sense:  $\Delta E < 0$  for the photon.

So  $|\Delta E| = \frac{2h^2 c^2}{m_e c^2 \lambda_0^2}$ , and we know  $|\Delta E| = 30 \text{ keV}$ .

$$\begin{aligned}
 \text{So } \lambda_0 &= \left( \frac{2(hc)^2}{m_e c^2 |\Delta E|} \right)^{1/2} \\
 &= \left( \frac{2(1240 \text{ eV}\cdot\text{nm})^2}{511 \text{ keV} \cdot 30 \text{ keV}} \right)^{1/2} = \boxed{1.42 \times 10^{-2} \text{ nm}}
 \end{aligned}$$

Is this consistent with  $\frac{2h}{m_e c} \ll \lambda_0$ ?

$$\text{Well, } \frac{2h}{m_e c} = \frac{2hc}{m_e c^2} = \frac{2 \cdot 1240}{511 \times 10^3} \text{ nm} = 4.85 \times 10^{-3} \text{ nm}$$

so  $\left( \frac{2h}{m_e c} \right) / \lambda_0 = 0.34$  which is kind of small.

Not a great approximation, but a decent one.

34 | From #30, we know  $\lambda_0 = \left( \frac{2(hc)^2}{m_e c^2 |\Delta E|} \right)^{1/2}$

$$\text{So } \lambda_0 = \left( \frac{2(1240 \text{ eV}\cdot\text{nm})^2}{511 \text{ keV} \cdot 50 \text{ keV}} \right)^{1/2} = 1.1 \times 10^{-2} \text{ nm}$$

$$\Rightarrow E = \frac{hc}{\lambda_0} = \frac{1240 \text{ eV}\cdot\text{nm}}{1.1 \times 10^{-2} \text{ nm}} = \boxed{113 \text{ keV}}$$

$$36] \quad \cancel{E_{\text{initial, photon}}} \quad E_{\text{final, photon}} = 80 \text{ keV}$$

$$E_e = 25 \text{ keV}$$

So by energy conservation,  $E_{\text{initial, photon}} = 105 \text{ keV}$

$$\Rightarrow \lambda = \frac{hc}{E} = \boxed{1.18 \times 10^{-2} \text{ nm}}$$

- 3-43 (a) A  $4000 \text{ \AA}$  wavelength photon is backscattered,  $\theta = \pi$  by an electron. The energy transferred to the electron is determined by using the Compton scattering formula  $\lambda' - \lambda_0 = \left(\frac{hc}{E_e}\right)(1 - \cos\theta)$  where we take  $E_e = m_e c^2$  for the rest energy of the electron

and so  $E_e \approx 0.511 \text{ MeV}$ . Upon substitution, one obtains

$$\Delta\lambda = 2(0.00243 \text{ nm}) = 0.00486 \text{ nm}.$$

The energy of a photon is related to its wavelength by the relation  $E = \frac{hc}{\lambda}$ , so the change in energy associated with a corresponding change in wavelength is given by  $\Delta E = -\left(\frac{hc}{\lambda^2}\right)\Delta\lambda$ .

Upon making substitutions one obtains the magnitude  $\Delta E = 6.0379 \times 10^{-24} \text{ J}$  and using the conversion factor 1 Joule of energy is equivalent to  $1.602 \times 10^{-19} \text{ eV}$ . The result is  $\Delta E = 3.77 \times 10^{-5} \text{ eV}$ .

- (b) This may be compared to the energy that would be acquired by an electron in the photoelectric effect process. Here again the energy of a photon of wavelength  $\lambda$  is given by  $E = \frac{hc}{\lambda}$ . With  $\lambda = 400 \text{ nm}$ , one obtains

$$E = \frac{(6.626 \times 10^{-34} \text{ J}\cdot\text{s})(3.0 \times 10^8 \text{ m/s})}{400 \times 10^{-9} \text{ m}} = 4.97 \times 10^{-19} \text{ J}$$

and upon converting to electron volts,  $E = 3.10 \text{ eV}$ .  $\frac{\Delta E}{E_{\text{photon}}} \approx 10^{-5}$ .

The maximum energy transfer is about five orders of magnitude smaller than the energy necessary for the photoelectric effect.

- (c) Could “a violet photon” eject an electron from a metal by Compton scattering? The answer is no, because the maximum energy transfer occurring at  $\theta = \pi$  is not sufficient.