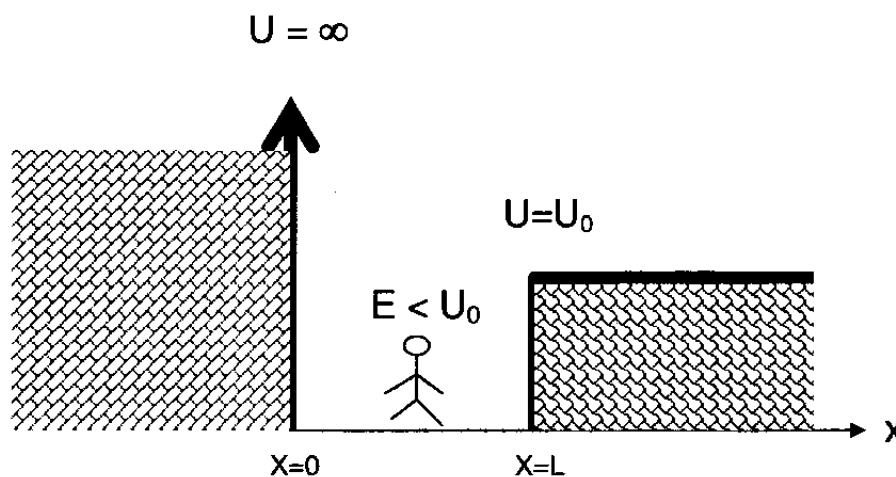


**Problem 1 : Drawing Lesson ! [12 pts] :** Consider a particle moving in a one-dimensional box with rigid walls ( $U=\infty$ ) at  $x = -L/4$  and  $x=L/4$ . (a) Sketch the potential in all regions of  $x$  , (b) write the wave functions and probability densities for the states  $n=1, 2, 3$  (c ) sketch the wave function and the probability density in all regions of  $x$ , (d) Calculate the ground state ( $n=1$ ) energy of the particle.

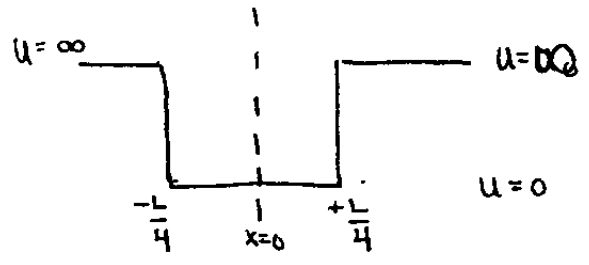
**Problem 2: Not Just Another Brick in the Wall ! [8 pts]:** Suppose that a particle of mass  $m$  is trapped in a potential well that has rigid wall at  $x = 0$  ( $U = \infty$  for  $x < 0$ ) and a finite wall of height  $U=U_0$  at  $x = L$ . See figure below and assume particle energy  $E < U_0$ . (a) Cleanly sketch the wave function for the lowest three states and be sure to label points  $x=0$  &  $x=L$  on your sketches, (b) What is the mathematical form of the wave function in the ground state in the three regions  $x < 0$ ,  $0 < x < L$  and  $X > L$  ?



# QUIZ 7

(i)

(a)



(b)

wavefunction must vanish for  $x > L/4$   
and  $x < -L/4$ , and the wavefunction  
must also be continuous everywhere, so  
 $\psi(+L/4) = \psi(-L/4) = 0$

solve Schrödinger's eqn:

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + U \psi = E \psi$$

inside the well,  $U=0$

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi$$

$$\frac{d^2 \psi}{dx^2} = -\frac{2mE}{\hbar^2} \psi = -k^2 \psi, \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

and the solutions to the differential eqn are

$$\psi(x) = A \sin(kx + \phi) + B \cos(kx + \phi)$$

let's impose the conditions  $\psi(-L/4) = \psi(+L/4) = 0$

$$\psi(-L/4) = A \sin\left(-\frac{kL}{4} + \phi\right) + B \cos\left(-\frac{kL}{4} + \phi\right) = 0$$

let's choose  $\phi = +\frac{kL}{4}$ , then

$$\begin{aligned} \psi(-L/4) &= A \sin\left(-\frac{kL}{4} + \frac{kL}{4}\right) + B \cos\left(-\frac{kL}{4} + \frac{kL}{4}\right) \\ &= 0 \end{aligned}$$

since  $\sin(0) = 0$ ,  $\cos(0) = 1$

we have

$$B = 0$$

$$\text{so } \psi(x) = A \sin\left(kx + \frac{kL}{4}\right)$$

now  $\psi(L/4) = 0 = A \sin\left(\frac{kL}{4} + \frac{kL}{4}\right) = A \sin\left(\frac{kL}{2}\right)$

for sin to vanish, we must have

$$\frac{kL}{2} = n\pi \Rightarrow k = \frac{2n\pi}{L}$$

$$\Rightarrow \psi(x) = A \sin\left(kx + \frac{kL}{4}\right)$$

$$= A \sin\left[k\left(x + \frac{L}{4}\right)\right]$$

$$\boxed{\psi_n(x) = A \sin\left[\frac{2n\pi}{L}\left(x + \frac{L}{4}\right)\right]}$$

(you could have also arrived at  $(x - L/4)$  depending upon the order in which you imposed the boundary conditions).

Now let's find the normalization constant:

$$A^2 \int_{-L/4}^{+L/4} dx \sin^2\left[\frac{2n\pi}{L}\left(x + \frac{L}{4}\right)\right] = 1$$

$$\text{let } u = \frac{2n\pi}{L}\left(x + \frac{L}{4}\right)$$

$$du = \frac{2n\pi}{L} dx$$

$$\text{so } dx = \frac{L}{2n\pi} du$$

limits change to

$$\int_0^{n\pi}$$

$$\frac{LA^2}{2n\pi} \int_0^{n\pi} du \sin^2 u = \frac{LA^2}{2n\pi} \left[ \frac{1}{2}u - \frac{1}{4}\sin 2u \right] \Big|_0^{n\pi}$$

$$= \frac{LA^2}{4} = 1$$

$$\Rightarrow A = \sqrt{\frac{4}{L}}$$

$$\boxed{\psi_n(x) = \sqrt{\frac{4}{L}} \sin\left[\frac{2n\pi}{L}\left(x + \frac{L}{4}\right)\right]}$$

$$\sin(A \pm B) = \sin A \cos B \pm \cos A \sin B$$

$$\Rightarrow \psi_n(x) = \sqrt{\frac{4}{L}} \sin \left[ \frac{2n\pi x}{L} + \frac{n\pi}{2} \right]$$

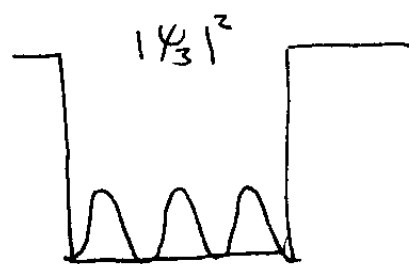
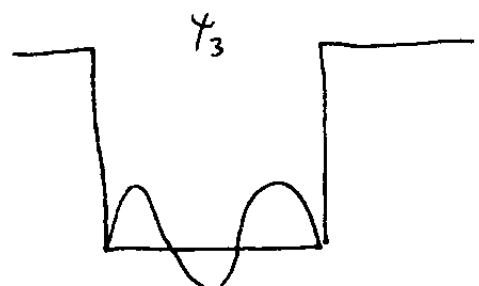
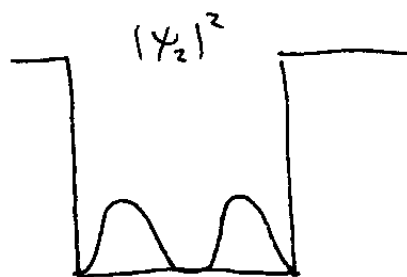
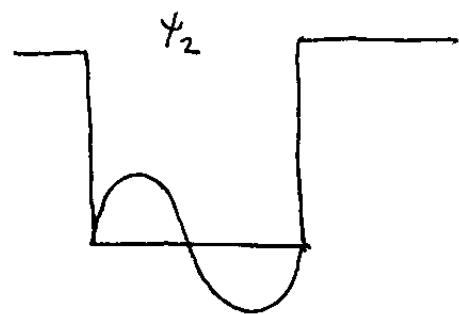
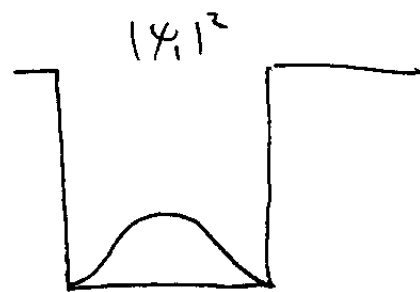
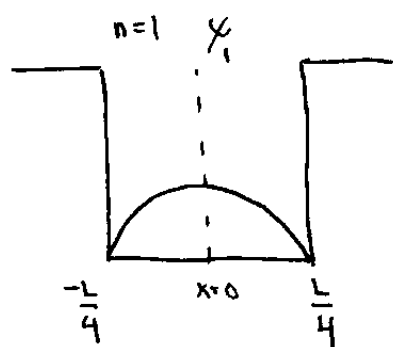
$$\psi_n(x) = \sqrt{\frac{4}{L}} \left[ \sin \left( \frac{2n\pi x}{L} \right) \cos \left( \frac{n\pi}{2} \right) + \cos \left( \frac{2n\pi x}{L} \right) \sin \left( \frac{n\pi}{2} \right) \right]$$

$$\text{for } n=1 \quad \psi_1(x) = \sqrt{\frac{4}{L}} \cos \left( \frac{2\pi x}{L} \right)$$

$$n=2 \quad \psi_2(x) = \sqrt{\frac{4}{L}} \sin \left( \frac{4\pi x}{L} \right) \quad (\text{forget about minus sign})$$

$$n=3 \quad \psi_3(x) = \sqrt{\frac{4}{L}} \cos \left( \frac{6\pi x}{L} \right) \quad (\text{forget about minus sign again})$$

the probability densities are just  $|\psi_n|^2$



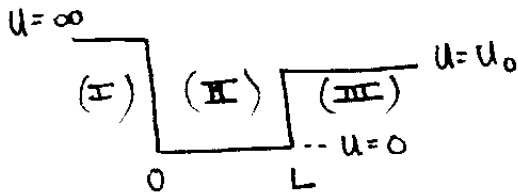
(d) the width of the well is  $\frac{L}{2}$  so the net energy is just

$$E_1 = \frac{\pi^2 \hbar^2}{2m \left(\frac{L}{2}\right)^2} = \frac{2\pi^2 \hbar^2}{mL^2}$$

(2)

$$\boxed{\psi^{(\text{I})}(x) = 0}$$

the wavefunction must vanish in region (I) since the potential is infinite here, so we must satisfy the condition  $\psi(0) = 0$  for the wavefunction to be continuous.



in region (II), Schrödinger's eqn reads

$$(0 < x < L) \quad -\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi \quad \text{since } U=0 \text{ here}$$

$$\frac{d^2 \psi}{dx^2} = -\frac{2mE}{\hbar^2} \psi = -k^2 \psi$$

$$\Rightarrow \psi^{(\text{II})}(x) = A \sin(kx + \phi) + B \cos(kx + \phi)$$

impose condition  $\psi^{(\text{II})}(0) = 0$

$$\Rightarrow \psi^{(\text{II})}(0) = A \sin \phi + B \cos \phi = 0$$

choose  $\phi = 0$ , and we obtain

$$B = 0$$

$$\Rightarrow \boxed{\psi^{(\text{II})}(x) = A \sin kx}$$

in region (III), Schrödinger's eqn is

$$(x > L) \quad -\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + U_0 \psi = E \psi, \quad U = U_0$$

$$\frac{d^2 \psi}{dx^2} = \frac{2m(U_0 - E)}{\hbar^2} \psi$$

$$\text{now } E < U_0 \text{ so } \frac{2m(U_0 - E)}{\hbar^2} > 0$$

$$\Rightarrow \frac{d^2 \psi}{dx^2} = \alpha^2 \psi \quad \alpha = \sqrt{\frac{2m(U_0 - E)}{\hbar^2}}$$

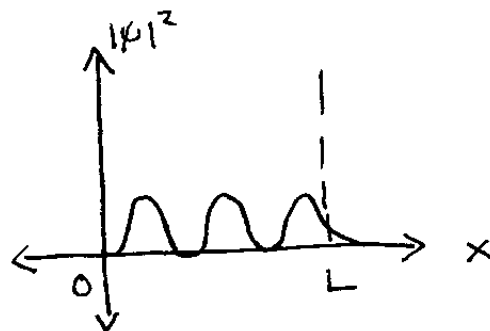
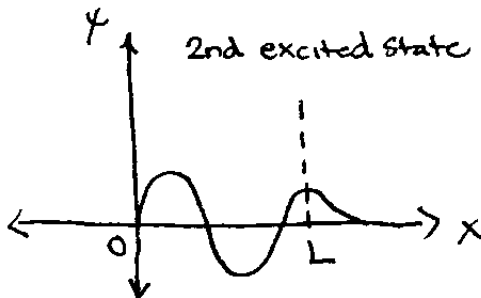
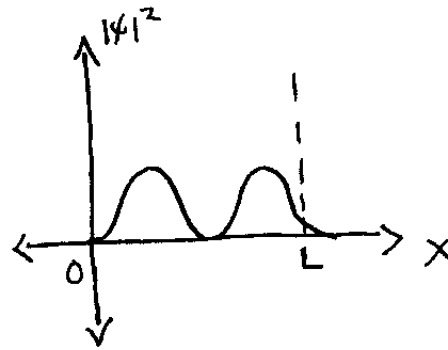
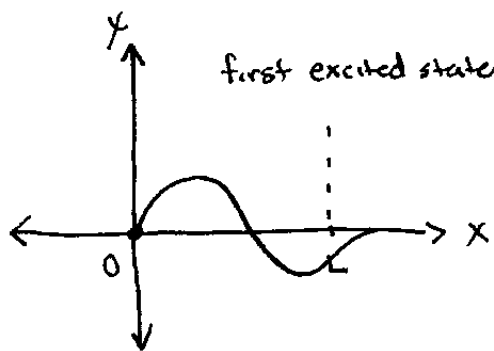
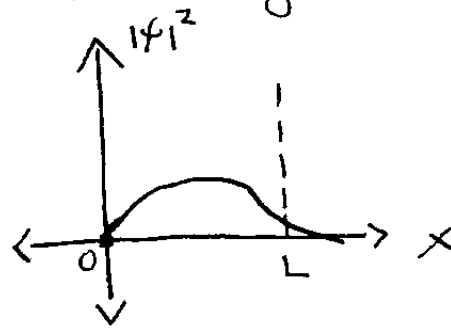
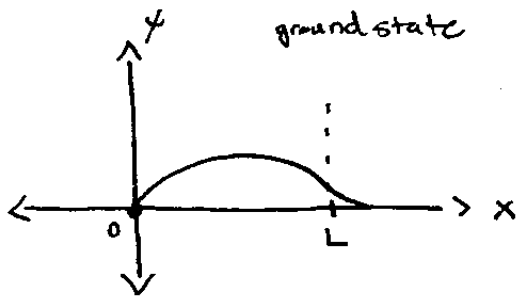
this has solutions  $\psi^{(III)}(x) = Ae^{\alpha x} + Be^{-\alpha x}$

but as  $x \rightarrow \infty$   $e^{\alpha x} \rightarrow \infty$ , so we set  $A=0$ , since the wavefunction must be finite everywhere.

$$\Rightarrow \boxed{\psi^{(III)}(x) = Be^{-\alpha x}}$$

so the wavefunctions and their probability densities look like

you did not have to sketch  $|\psi|^2$



~~you did not have to sketch |ψ|^2~~